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Acoustics of Excited Jets—A Historical Perspective

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Abstract

The idea that a jet may be excited by external forcing is not new. The first published demonstration of a jet responding to external pressure waves occurred in the mid-1800's. It was not, however, until the 1950's, with the advent of commercial jet aircraft, that interest in the subject greatly increased. Researchers first used excited jets to study the structure of the jet and attempt to determine the nature of the noise sources. The jet actuators of the time limited the range (Reynolds and Mach numbers) of jets that could be excited. As the actuators improved, more realistic jets could be studied. This has led to a better understanding of how jet excitation may be used not only as a research tool to understand the flow properties and noise generation process, but also as a method to control jet noise.

1. Introduction

The concept that a fluid flow may be influenced by external sources has its origins in the 19th century observations of Leconte (ref. 1), Tyndall (ref. 2), and, later, Rayleigh (ref. 3). However, it was not until the advent of commercial jet engines in the 1950's that others began to further explore these ideas in search of quieter aircrafts and a better understanding of the fluid mechanics in the jet.

The subject of flow excitation covers a vast array of research. As such, the paper presented here is not an exhaustive study of all excited flow research conducted. Rather, this paper seeks to give an overview of the physics of excited jets, primarily in terms of knowledge required to excite a jet. A brief discussion of flow actuators, whose development is crucial if research is to expand to higher Reynolds number jets outside of the laboratory environment, is then addressed. Finally, two research topics, in which jet excitation has played an important role, will also be discussed. The first topic is the application of jet excitation to influence and understand the large-scale coherent structures found in nearly all turbulent flows. While these structures do not directly produce noise (in a subsonic jet), they govern the breakdown of the flow into small-scale turbulence that produces jet noise. Thus, control, exerted in the form of excitation, of the large-scale structures leads to changes in the mixing properties of the jet and, ultimately, the sound produced. The second research topic deals with the direct implications of excitation on the radiated sound. Jets, excited under certain conditions, are observed to produce more broadband noise than in their unexcited state. Excited under other conditions, however, they are observed to produce less broadband noise than in their unexcited state. Following these observations, considerable research has been conducted to determine the mechanism coupling the excitation conditions to the noise produced. The research shows that jet excitation can alter primary properties (turbulent structure, plume mixing rate, and noise produced) of a turbulent shear flow. With a greater understanding of the physics at work and advanced flow, the potential exists for controlled excitation to optimize the properties of a jet for any desired application.

2. Basic Principles of Jet Excitation

Jet excitation occurs when a perturbation alters the instabilities already present in the shear layer between the jet and the ambient medium. These perturbations may be created naturally or by intentionally applied forces. While engineers may have little control over the natural perturbations, there are several

reasons they may want to introduce artificial excitation. As a research tool, jet excitation is used to study flow instabilities, often in pursuit of enhanced mixing for chemical processes, heat transfer, or plume reduction (ref. 4). Artificial excitation has been used to study large-scale coherent structures within the jet plume. In this application, the excitation amplifies the structures of interest and locks them into a regular frequency allowing use of more conventional measurement techniques (ref. 4). Finally, artificial excitation may suppress the turbulence in a free shear flow.

As there are several reasons to introduce excitation to a jet the specific application dictates which instabilities are targeted for excitation. The boundary layer between two fluids with different velocities, temperatures, and densities is inherently unstable. As the fluids mix in search of equilibrium, the boundary layer expands forming Kelvin-Helmholtz instability waves. As the instability waves grow, they break down and dissipate leaving the fluid fully mixed in equilibrium. While the instability waves do not directly radiate sound in a subsonic jet, they govern the formation of the turbulence that does create sound (ref. 5). In a supersonic jet, the instability waves themselves radiate significant sound (ref. 6).

The behavior of the instability waves may be decomposed into modes of oscillation for convenience of analysis. Each mode is characterized by an initial growth (or amplification) rate, governed by linear stability theory and a saturation limit determined by nonlinear effects (ref. 7). Additionally, each mode has a preferred frequency (usually written in nondimensional form as Strouhal number) where maximum amplitude may be obtained. Higher frequency components grow faster and, therefore, saturate closer to the nozzle exit than the lower frequency components that persist further downstream (ref. 7). Within this framework, it has been shown by Plaschko that the thin axisymmetric shear layer near the nozzle exit is unstable to a large number of discrete azimuthal modes (ref. 8). In the fully developed jet downstream of the potential core, however, Batchelor and Gill showed that the jet is only unstable to the first helical mode ($m = 1$) (ref. 9). Cohen and Wygnanski showed the axial evolution of the lowest seven azimuthal modes ($0 \leq m \leq 6$) by calculating the amplification factor at locations from the nozzle exit to four jet diameters downstream in a $\sim 2.7 \times 10^4$ Reynolds number jet (ref. 10). It was also observed by Corke, Shakib, and Nagib that while they have similar growth rates, the axisymmetric and first helical modes do not coexist at the same time and place (ref. 11).

Experimentally, many studies have investigated the growth and saturation of instability modes. In a systematic study of jet response to plane wave excitation at several frequencies (axisymmetric mode or $m = 0$), Crow and Champagne found that the mode with the maximum attainable amplitude (“preferred mode”) at the end of the potential core corresponded to the mode with Strouhal frequency 0.3 (ref. 7). In a later study, Gutmark and Ho showed that the Strouhal frequency of this “preferred mode” varied between experimental jet facilities and concluded that the specific boundary layer characteristics play an important role in the selection of the most amplified mode (ref. 12). Moore excited the axisymmetric mode using acoustic plane waves of different amplitudes. The results show that considerable effects on noise and mixing occur at low levels of excitation if the excitation frequency is near the preferred frequency of the jet (ref. 5). Moore concluded that many possible sources of excitation exist at this level in an actual engine, such as combustion noise. Others have investigated instability growth and saturation in the lowest order helical modes ($m \leq 2$) as well as the axisymmetric mode: Chan ($m = 0, 1, 2$) (ref. 13), Strange and Crighton ($m = 0, 1, 2$) (ref. 14), and Cohen and Wygnanski (ref. 10).

Acoustically, only a small subset of possible instability modes contributes to the sound produced for low Mach number jets. In 1975, Michalke and Fuchs applied Lighthill’s analogy to an axisymmetric jet and showed that only the lowest three or four azimuthal modes significantly contribute to the sound field (ref. 15). In later theoretical work, Kopiev and Chernyshev found a similar result by studying vortex ring eigen-oscillations (ref. 16). This theoretical work is supported by experimental observations by Armstrong et al. in the acoustic near field (ref. 17) and by Juve et al. in the acoustic far field (ref. 18).

Most excited jet research has been done in the laboratory on small cold jets with relatively low initial turbulence levels and at relatively low Reynolds number conditions compared to actual engine applications which are hot and have very high Reynolds numbers. A few researchers have investigated the effects of heat and initial turbulence on jet excitability. Lepicovsky, Ahuja, and Salikuddin studied tone-excited heated jets and concluded that the jet sensitivity to excitation is a strong function of jet velocity

(ref. 19). In general, the researchers found that the addition of heat decreased the excitation level required for a low Mach number jet ($M = 0.3$) to respond while increasing the excitation level needed for a higher Mach number jet to respond ($M = 0.8$). However, they found no significant change in preferred excitation Strouhal number (ref. 19). Others investigating excited heated jets include Jubelin (ref. 20), Lu (ref. 21), and Vermeulen, Odgers, and Ramesh (ref. 22).

The effect of initial turbulence intensity on jet response to excitation has also been investigated. Raman, Zaman and Rice performed a series of experiments in 1989 where the initial turbulence was varied from 0.15 to 5 percent using screens and boundary layer trips placed inside the jet (ref. 23). The results show that the jet at all initial turbulence levels responded to the excitation; the amplitude required to reach the same level of response, however, increased as the initial turbulence intensity increased. It was also reported that the preferred excitation frequency varied little with changes in turbulence intensity. A similar study, Mankbadi, Raman and Rice, varied only the core turbulence intensity while holding the turbulence intensity in the boundary layer as steady as possible (ref. 24). The researchers found that the jet response was reduced as core turbulence intensity increased. Thus, more turbulent flows require more powerful actuators to excite them. This is a critical issue in developing actuators for practical applications.

3. Flow Actuators

Flow actuators generate forces that act on the flow and create the excitation effect. Depending on the actuator, the forces generated may assume different forms. The three common excitation forces are mechanical, pressure (acoustic) waves, and fluidic interactions. Each actuation method carries advantages and disadvantages. Generally, the power, frequency response, and number of actuators must be balanced against the size, weight, and expense of the actuators when selecting an excitation system for a particular application. The following section will discuss physical, pressure, and fluidic actuation methods.

A. Physical Actuators

Perhaps the simplest excitation method, physical actuators periodically place solid objects in the flow to influence its behavior. This has the effect of redirecting the flow and changing the amplitude of the instability waves. These physical barriers have the advantage of being very strong, able to exert large forces capable of completely redirecting the flow. Frequency response of the actuators, however, is often limited.

A classic example of physical excitation is the use of vibrating ribbons to study Tollmein-Schlichting waves generated during boundary layer transition (ref. 4). Following this idea, Oster and Wygnanski used a thin flap, rotated about one end and driven by two voice coils, where used to study turbulent mixing between two excited streams (ref. 25). Along similar lines, very low Reynolds number bifurcating and blooming jets can be generated by “wobbling the nozzle slightly” (ref. 26).

A modern mechanical actuator combines the forces of physical and fluidic forces. Piezoelectric actuators, like those developed by Wiltse and Glezer (ref. 27), move a small wedge a small distance into the flow. The wedge, which oscillates into and out of the flow, creates fluidic vortices in its wake. These vortices generate the actual excitation force. Piezoelectric actuators have been used to study mixing enhancement (refs. 27 and 28) and to excite the smallest scales, within the turbulence dissipation range, in a free shear flow (ref. 29).

B. Acoustic Actuators

Acoustic waves were the exciting force for the earliest observations of excited flows. In fact, most experimental research on excited flows employed acoustic waves, generated by loudspeakers, as the exciting force. Loudspeakers offer several advantages for use in excited jet research. They are fairly inexpensive and robust. Loudspeakers allow simple control of frequency, relative phase, and amplitude.

While this makes them a convenient tool, there are limitations on the maximum wave amplitude and frequency that may be generated, restricting the size and speed of the jet that may be influenced. In addition, their size and weight make loudspeakers a poor choice for any application that requires mobility (real-life aircraft engines for example). As rapid mobility is generally not required in a laboratory using small model scale jets, loudspeakers are a good choice. A single loudspeaker, placed in the plenum chamber of a jet rig, is the most common method used for exciting the axisymmetric mode of the jet. Several loudspeakers, mounted at the nozzle exit and set to an appropriate relative phasing, are frequently used to excite the higher modes of the jet.

C. Fluidic Actuators

Fluidic actuators assume many forms. From the very simple to the most complex, fluidic actuators function by forcing fluid into the jet plume in the form of pulses or vortices. A simple example, presented by Brown and Ahuja in a 1990 proof of concept paper (ref. 30), consists of a vortex shedding ring placed shortly upstream of the nozzle exit where the fan stream or free stream fluid could generate the exciting vortices. The frequency of the vortex generation is controlled by the diameter of the rod used to build the shedding ring.

Another actuator, offering more control than a vortex shedding ring, is the “flip-flop” jet or fluidic actuator. The flip-flop jet uses no moving parts but rather relies on the fluid itself to switch paths according to the naturally varying pressure gradient (ref. 31). A feedback tube, whose length and controls the frequency of oscillation (ref. 32), creates regularity in the switching. Flip-flop jets are capable of operating at supersonic exit conditions creating the potential for very large amplitude perturbations (ref. 33). The frequency of oscillation, however, may be a limiting factor (ref. 33). Excitation using flip-flop nozzles has been evaluated for use in jet mixing enhancement (ref. 34).

Synthetic, or zero mass jets offer another fluidic excitation device. By varying the pressure in a small orifice, fluid from the flow itself can be periodically sucked in and expelled creating the fluidic pulses that excite the jet (ref. 35). The pressure within the orifice can be controlled using speaker type membranes or by piston actuation. Piston driven synthetic jets have the advantages of higher operating pressures and broader frequency range (ref. 36). Synthetic jets are unique because the net mass flux of the jet is zero.

Spark driven actuators form a final class of fluidic actuators. In this class the exciting force is created by a rapid expansion of fluid due to a combustion process or due to temperature expansion. Combustion driven actuators, investigated by Crittenden, Glezer, Funk, and Parekh in 2001 (ref. 36), burn a gaseous fuel/air mixture in a chamber, expelling the exhaust through an orifice where it acts to excite the flow. The combustion driven actuator has a frequency response limited by the total time required to fill the chamber and burn the contents (which is a function of the chamber size). Another spark driven actuator avoids the problems of combustion by placing a high temperature spark directly in the air. The plasma actuator, developed by Samimy et al. (ref. 37), creates a rapid expansion of the air around the nozzle generating the forcing effect. Plasma actuators have good frequency response and have been shown effective on jets up to Mach 1.3 (ref. 38).

Fluidic actuators have, like the flip-flop jet, synthetic jet, and plasma actuator, have developed because of the high amplitude, high frequency forcing required to excite a turbulent high Reynolds number jet. They now allow excitation of jets not possible with the traditional acoustic loudspeaker and, thus, further understanding into the physics of turbulent jets.

4. Applications

A. Study of Large-Scale Coherent Structures and Mixing Enhancement

Coherent structures are generated by flow instabilities and their formation is governed by the initial condition for the flow where the structure originates (ref. 39). While these structures do not directly

radiate sound, they govern the production of the small-scale turbulence that is responsible for generating sound (ref. 5). To maintain control over a turbulent flow, Crow and Champagne theorized that either the vorticity must be forced directly using body forcing or the boundary conditions must be controlled at all times (ref. 7). As it is not practical to directly force the vorticity everywhere in the flow, the only avenue of control remaining is the boundary conditions. But, if coherent structures are indeed governed by the initial condition, controlling the boundary condition at the nozzle exit is a logical method to influence at least the initial generation of coherent structures. For this reason, control of the jet seems to be limited to the initial region of the jet. Hussain and Zaman noted that no phase-locked coherent structures could be measured beyond 6 jet diameters downstream of the nozzle exit (ref. 40). Additionally, the turbulence intensity in the forced jet trends to the same levels as that in the unforced jet between six and eight jet diameters downstream of the nozzle exit (ref. 7).

Studies dealing with large-scale structures often apply excitation techniques. Although it is possible to measure large-scale structures without excitation ((ref. 41) for example), the addition of excitation serves to amplify the coherent structures above the background turbulence levels and lock them into a steady formation and pairing frequency. Under the influence of excitation, structures are repeatable in time and space allowing for more traditional phase-average measurements (ref. 42).

In addition to numerous studies on subsonic flows ((refs. 7, 42, 43, 44), others), large-scale coherent structures have also been studied using excitation methods in supersonic flows by Lepicovsky et al. (ref. 45). Using axisymmetric excitation in a fully expanded, shock free supersonic jet (Mach 1.37), the researchers found that the relationship between phase velocity and excitation Strouhal frequency follows the same general trends observed in the subsonic jets. In supersonic jets, these large-scale structures do radiate noise directly and, thus, excitation could play a direct role in altering the sound production.

Large-scale coherent structures in turbulent flows are important energy transport mechanism (ref. 39). Ho and Huang found that multiple vortices converge simultaneously, increased the rate of spreading in a mixing layer, under low-level excitation at a subharmonic frequency (ref. 43). Jet plume mixing enhancement, using a single frequency plane wave, has a maximum limit (ref. 46). This finding led to several studies of using excitation at multiple frequencies simultaneously. Higher spreading rates were found when the jet was excited at the fundamental and subharmonic frequencies (ref. 47). Higher spreading rates were also found when the jet was excited in both the axisymmetric and first helical mode simultaneously (ref. 48). Exerting control over the spreading rate significantly changes the production of small-scale turbulence, the noise generation mechanism in a subsonic jet.

B. Broadband Noise Amplification

In 1975, Bechert and Pfizenmaier presented some interesting acoustic measurements from a jet, excited by a pure tone, showing an amplification of the broadband noise of 6–7 dB (ref. 49). Moore, in 1977, confirmed the phenomenon and further showed that it can be created by a low level excitation (ref. 5). Deneuille and Jacques showed that amplification was also possible when a broadband excitation source was used (ref. 50). Bechert and Pfizenmaier, also in 1977, generated the broadband amplification effect by exciting the first and second azimuthal modes of the jet (ref. 51). In both cases (broadband source or higher mode excitation), the characteristic tones, which accompanied the broadband amplification in previous experiments, were absent. Jubelin considered both hot jets and supersonic jets and showed that both exhibit broadband amplification when excited (ref. 20). Then, in 1980, Kibens revealed an experiment showing a reduction in broadband noise by up to 10 dB in an excited jet (ref. 52). In 1985, researchers from the Lockheed-Georgia Company examined the jet noise broadband amplification problem using flow visualization (ref. 44), flow measurements (ref. 53) and acoustic measurements (ref. 54). They concluded that the amplified large-scale structures result in increased small-scale turbulence and that small-scale turbulence is responsible for the increased broadband noise (ref. 54). Zaman also examined the problem of broadband amplification and theorized that the forced vortex pairings created by the excitation were responsible for the increased noise (ref. 55). However, neither theory explains the noise reduction observed by Kibens.

Jet excitation can increase turbulence and enhance mixing. In 1981, Zaman and Hussain showed that certain excitation conditions would suppress turbulence intensities in jet with laminar initial boundary layers (ref. 56). Flow visualization led to the conclusion that the large-scale structures were destroyed by the excitation before they fully developed, reducing the large-scale fluctuations and, therefore, the turbulence intensity. In a follow on study, Hussain and Hasan reported that the jet noise was reduced when the turbulence was suppressed (ref. 57). These results were not repeatable when the initial boundary layer was turbulent. Thus, for very low Reynolds number jets with laminar boundary layers turbulence levels and, as a result broadband noise levels, can be suppressed. It should be noted that, unlike broadband noise amplification, reliable suppression of broadband noise has not been observed. This is due, in part, to the fact that most research uses excitation that enhances turbulence structures and increases mixing, both of which tend to produce increased noise.

5. Conclusions

Excitation is one of the very few ways to influence a jet after it passes the nozzle exit. The frequency and amplitude of the perturbations placed on the flow dictate, to a large extent, the behavior of the jet as it develops in the first few jet diameters downstream of the nozzle exit. While considerable work has been done in some areas, there are other areas where there is still much to explore. For example, little research exists on the excitation of azimuthal modes above two. Theory predicts that these modes do not efficiently radiate sound. Investigation, using excitation to amplify these modes, could confirm this and offer insights into the mechanisms at work.

The idea that external forcing at particular levels and frequencies can change the dynamics of a jet is powerful. Research has shown how excitation can enhance certain structures within low Reynolds number jets to dramatically change the turbulence and noise characteristics of the plume. Very little of this research, however, studies the highly turbulent, high Reynolds number jets likely found in real-world applications. This is primarily due to the high frequency and amplitude requirements needed to excite a highly turbulent flow. But, actuator technology is progressing and devices like the flip-flop jets, synthetic jets, and plasma actuators are significantly more capable of dealing with the higher requirements of turbulent flows than their predecessors were. Armed with these latest flow actuators and the understanding developed studying the low Reynolds number jets, the next development in excitation technology is poised to move it out of the laboratory and toward more everyday applications.

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